

QUADRUPLE SERIES EQUATIONS INVOLVING HEAT POLYNOMIALS

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ABSTRACT

In this paper, an exact solution of the quadruple series equations involving heat polynomials $P_{n,v}(x,t)$ is given. Certain integral and series representations for $P_{n,v}(x,t)$ and $w_{n,v}(x,t)$ are shown, which are needed in the course of analysis.

Introduction

In this paper, we consider the following quadruple series equations:

$$\sum_{n=0}^{\infty} A_n \frac{t^{-n} \rho^n}{\Gamma\left(v + \frac{1}{2} + n + p\right)} P_{n+p,v}(x,-t) = \phi_1(x,t); 0 \leq x < y \quad (1.1)$$

$$\sum_{n=0}^{\infty} A_n \frac{1}{\Gamma\left(\mu + \frac{1}{2} + n + p\right)} P_{n+p,\sigma}(x,-t) = \phi_2(x,t); y < x < z \quad (1.2)$$

$$\sum_{n=0}^{\infty} A_n \frac{t^{-n} \rho^n}{\Gamma\left(v + \frac{1}{2} + n + p\right)} P_{n+p,v}(x,-t) = \phi_3(x,t); z < x < h \quad (1.3)$$

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$$\sum_{n=0}^{\infty} A_n \frac{1}{\Gamma\left(\mu + \frac{1}{2} + n + p\right)} P_{n+p,\sigma}(x, -t) = \phi_4(x, t); h < x < \infty \quad (1.4)$$

Where, $\phi_1(x, t), \phi_2(x, t), \phi_3(x, t)$ and $\phi_4(x, t)$ are prescribed functions for $t > \rho > 0$ and A_n is to be determined, and $P_{n,v}(x, t)$ is the heat polynomial [Haimo, 1966] defined by

$$P_{n,v}(x, t) = \sum_{k=0}^n 2^{2k} \binom{n}{k} \frac{\Gamma\left(v + \frac{1}{2} + n\right)}{\Gamma\left(v + \frac{1}{2} + n - k\right)} x^{2n-2k} t^k, v > 0 \quad (1.5)$$

It may be noted that $P_{n,o}(x, t) = v_{2n}(x, t)$ is the ordinary heat polynomial of even order defined by Rosenbloom and Widder (1959) and that

$$P_{n,o}(x, -1) = (-1)^n \square 2^{2n} n! \square L_n^{-\frac{1}{2}}\left(\frac{x^2}{4}\right) = H_{2n}\left(\frac{x}{2}\right), \text{ the Hermite polynomial of even order}$$

defined by Erdelyi (1853). The analysis given here is purely formal and no attempt is made to supply details of rigorous proof.

2. SOLUTION

Multiplying equation (1.1) by $t^{-\left(p+m+\mu+\frac{1}{2}\right)} \square (t-\alpha)^{\mu+m-v-1}$ where m is a positive integer, integrating with respect to t from ρ to ∞ and using (2.11) we get

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu + m + n + p + \frac{1}{2}\right)} P_{n+p,\mu+m}(\xi, -\rho) = F_1(\xi, \rho), 0 < \xi < y \quad (3.1)$$

Where $\mu + m > v - \frac{1}{2}$ and

$$F_1(\xi, \rho) = \frac{\rho^{\nu+\frac{1}{2}}}{\Gamma(\mu+m-\nu)} \int_{\rho}^{\infty} t^{-(\rho+m+\mu+\frac{1}{2})} \square (t-\rho)^{\mu+m-\nu-1} \phi_1(\xi, t) dt \tag{3.2}$$

Further setting $t = \rho$ in (1.2), multiplying it by $x(x^2 - \xi^2)^{\sigma-\mu-1} \square e^{-\frac{x^2}{4}}$, integrating with respect to x from ξ to ∞ and applying (2.13) we obtain

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu+n+\rho+\frac{1}{2}\right)} P_{n+\rho, \mu}(\xi, -\rho) = F_2(\xi, \rho), y < \xi < z \tag{3.3}$$

Where $\sigma > \mu > -\frac{1}{2}$ and

$$F_2(\xi, \rho) = 2^{1-2(\sigma-\mu)} \square e^{-(\sigma-\mu)} \square e^{\left(\frac{\xi^2}{4\rho}\right)} \int_{\xi}^{\infty} x(x^2 - \xi^2)^{\sigma-\mu-1} \square e^{-\left(\frac{x^2}{4\rho}\right)} \square \phi_2(x, \rho) dx \tag{3.4}$$

Now multiplying (3.1) by $\xi^{2(\mu+m)-1}$ and applying the operator $\left(\frac{d^2}{d\xi^2}\right)^m$, we see in view of (2.6) that

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu+\frac{1}{2}+n+\rho\right)} P_{n+\rho, \mu}(\xi, -\rho) = H(\xi, \rho), 0 < \xi < y \tag{3.5}$$

Where $H_1(\xi, \rho) = \xi^{1-2\mu} \left(\frac{d^2}{d\xi^2}\right)^m \left[\xi^{2(\mu+m)-1} F_1(\xi, \rho)\right]$ (3.6)

Now, we multiply equation (1.3) by $t^{-(\rho+m+\mu+\frac{1}{2})} \square (t-\alpha)^{\mu+m-\nu-1}$, where m is a positive integer, integrating with respect to t form ρ to ∞ and using (2.11), we get

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu+m+n+\rho+\frac{1}{2}\right)} P_{n+\rho+\mu+m}(\xi, -\rho) = F_3(\xi, \rho), z < \xi < h \tag{3.7}$$

Where, $\mu+m > \nu > -\frac{1}{2}$ and

$$F_3(\xi, \rho) = \frac{\rho^{\nu+\frac{1}{2}}}{\Gamma(\mu+m-\nu)} \int_{\rho}^{\infty} t^{-(\rho+m+\mu+\frac{1}{2})} \square (t-\rho)^{\mu+m-\nu-1} \phi_3(\xi, t) dt \tag{3.8}$$

Setting $t = \rho$ in (1.4), multiplying it by $x(x^2 - \xi^2)^{\sigma-\mu-1} \square e^{-\frac{x^2}{4}}$, integrating with

respect to x from ξ to ∞ and applying (2.13) we obtain.

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu + n + \rho + \frac{1}{2}\right)} P_{n+\rho, \mu}(\xi, -\rho) = F_4(\xi, -\rho), h < \xi < \infty \quad (3.9)$$

Where $\sigma > m > -\frac{1}{2}$ and

$$F_4(\xi, \rho) = 2^{1-2(\sigma-\mu)} \rho^{-(\sigma-\mu)} \frac{e^{\left(\frac{\xi^2}{4\rho}\right)}}{\Gamma(\sigma-\mu)} \int_{\xi}^{\infty} x(x^2 - \xi^2)^{\sigma-\mu-1} e^{-\left(\frac{x^2}{4\rho}\right)} \phi_4(x, \xi) dx \quad (3.10)$$

Now, multiplying (3.7) by $\xi^{2(\mu+m)-1}$ and applying the operator $\left(\frac{d}{d\xi^2}\right)^m$, we see in view of (2.6) that

$$\sum_{n=0}^{\infty} \frac{A_n}{\Gamma\left(\mu + \frac{1}{2} + n + \rho\right)} P_{n+\rho, \mu}(\xi, -\rho) = H_2(\xi, -\rho), z < \xi < h \quad (3.11)$$

$$\text{Where } H_2(\xi, \rho) = \xi^{1-2\nu} \left(\frac{d^2}{d\xi^2}\right)^m \left[\xi^{2(\mu+m)-1} \phi_3(\xi, \rho)\right] \quad (3.12)$$

The left hand sides of (3.11), (3.5), (1.2), (1.4) are now identical and an application of the orthogonality relation (2.2) gives the solution of the quadruple series equations (1.1) - (1.4) in the form

$$A_n = \frac{\Gamma\left(\mu + \frac{1}{2}\right)}{2^{4(n+\rho)} (n+\rho)!} \left[\int_0^y W_{n+\rho, \mu}(\xi, \rho) H_1(\xi, \rho) dR(\xi) + \int_z^z W_{n+\rho, \mu}(\xi, \rho) \phi_2(\xi, \rho) dR(\xi) + \int_h^y W_{n+\rho, \mu}(\xi, \rho) \phi_3(\xi, \rho) dR(\xi) + \int_h^z W_{n+\rho, \mu}(\xi, \rho) H_2(\xi, \rho) dR(\xi) \right] \quad (3.13)$$

Where $H_1(\xi, \rho)$ and $H_2(\xi, \rho)$ are the same as defined by (3.6) and (3.12) respectively and $dR(\xi)$ is defined by (2.3).

Using the relation (2.1) and setting $B_n = A_n (-1)^{n+p} 2^{2(n+p)} (n+p)!$, we find that the equations (1.1), (1.2), (1.3) and (1.4) transforms into

$$\sum_{n=0}^{\infty} \frac{B_n \rho^n}{\Gamma\left(\nu + \frac{1}{2} + n + p\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{x^2}{4t}\right) = t^{-p} \phi_1(x, t), 0 < x < y \tag{3.14}$$

$$\sum_{n=0}^{\infty} \frac{B_n t^n}{\Gamma\left(\mu + \frac{1}{2} + n + p\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{x^2}{4t}\right) = t^{-p} \phi_2(x, t), y < x < z \tag{3.15}$$

$$\sum_{n=0}^{\infty} \frac{B_n \rho^n}{\Gamma\left(\nu + \frac{1}{2} + n + p\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{x^2}{4t}\right) = t^{-p} \phi_3(x, t), z < x < h \tag{3.16}$$

$$\sum_{n=0}^{\infty} \frac{B_n t^n}{\Gamma\left(\mu + \frac{1}{2} + n + p\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{x^2}{4t}\right) = t^{-p} \phi_4(x, t), h < x < \infty \tag{3.17}$$

Where $t \geq \rho > 0$ and their solution is given by

$$B_n = \frac{\Gamma\left(\mu + \frac{1}{2}\right) \Gamma(n+p+1)}{2^{\nu + \frac{1}{2}} \rho^{\nu + \frac{1}{2} + 2n}} \left[\begin{aligned} & \int_0^y e^{-\left(\frac{\xi^2}{4\rho}\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{\xi^2}{4\rho}\right) H_1(\xi, \rho) dR(\xi) \\ & + \int_y^z e^{-\left(\frac{\xi^2}{4\rho}\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{\xi^2}{4\rho}\right) \phi_2(\xi, \rho) dR(\xi) \\ & + \int_z^h e^{-\left(\frac{\xi^2}{4\rho}\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{\xi^2}{4\rho}\right) \phi_3(\xi, \rho) dR(\xi) \\ & + \int_h^{\infty} e^{-\left(\frac{\xi^2}{4\rho}\right)} L_{n+p}^{\left(\nu - \frac{1}{2}\right)}\left(\frac{\xi^2}{4\rho}\right) H_2(\xi, \rho) dR(\xi) \end{aligned} \right] \tag{3.18}$$

The solution of quadruple series equations involving generalized Laguerre polynomials can be obtained independently by the above procedure.

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